42. Method of Descent

In this exercise we will apply the method of descent to solve the wave equation on \mathbb{R}^2 for a particular set of initial conditions. The idea is to help you understand the key ideas and notation of the method. It is a combination of results from Sections 5.1–5.

Consider the wave equation on \mathbb{R}^2 with initial conditions

$$\partial_t^2 u - \Delta u = 0 \text{ on } (x, t) \in \mathbb{R}^2 \times (0, \infty),$$

 $u(x, 0) = g(x) = \chi_{[0, \infty)}(x_1), \qquad \partial_t u(x, 0) = h(x) = 0.$

- (a) Suppose u is a solution of the wave equation for n=2. Why does $(x_1, x_2, x_3, t) \mapsto u(x_1, x_2, t)$ solve the wave equation on \mathbb{R}^3 ? (note, the Laplacians are different in different dimensions).
- (b) Conversely, prove that a solution \bar{u} to the 3-dimensional wave equation that does not depend on x_3 gives a solution to the 2-dimensional wave equation. (1 point)
- (c) By (a) and (b), we now must solve a wave equation on \mathbb{R}^3 . The key to solving the 3-dimensional wave equation is to consider the (spatial-)spherical means

$$U(x,t,r) = \frac{1}{4\pi r^2} \int_{\partial B(x,r)} \bar{u}(z,t) \; d\sigma(z),$$

and likewise let G and H be the spherical means of \bar{g} and \bar{h} respectively. Explain why $\bar{g}(x_1, x_2, x_3) = \chi_{[0,\infty)}(x_1)$ and $\bar{h} = 0$ (or give the definition of bar). Show that

$$G(x,r) = \begin{cases} 0 & \text{for } x_1 \le -r \\ \frac{1}{2} \frac{x_1 + r}{r} & \text{for } |x_1| \le r \\ 1 & \text{for } r \le x_1 \end{cases} \text{ and } H(x,r) = 0.$$

You may use the following geometric fact: for -R < a < b < R, the surface area of the part of the sphere $\partial B(0,R)$ with $a < x_1 < b$ is $2\pi R(b-a)$. (4 points)

(d) We know by Lemma 5.2 that U obeys the Euler-Poisson-Darboux equation. Let $\tilde{U}(x,t,r):=rU(x,t,r)$. Show that \tilde{U} obeys the following PDE

$$\begin{split} \partial_t^2 \tilde{U} - \partial_r^2 \tilde{U} &= 0 \text{ on } (t,r) \in [0,\infty) \times [0,\infty), \\ \tilde{U}(x,0,r) &= rG(x,r), \qquad \partial_t \tilde{U}(x,0,r) = rH(x,r). \end{split}$$

Note that there are no x-derivatives in this PDE, so we can think of it as a family of PDEs in the variables r, t parametrised by x. (2 points)

(e) Thus we see that \tilde{U} obeys the 1-dimensional wave equation on the half-line $r \in [0, \infty)$. This is solved by a trick using reflection, and the formula is at the end of Section 5.2. We only need the solution for small r, so it is enough to consider the case $0 \le r \le t$. In this case,

show

$$\tilde{U}(x,t,r) = \begin{cases} 0 & \text{for } x_1 \le -(t+r) \\ \frac{1}{4}(x_1+t+r) & \text{for } -(t+r) \le x_1 \le -(t-r) \\ \frac{1}{2}r & \text{for } |x_1| \le t-r \\ \frac{1}{4}(x_1-t+3r) & \text{for } t-r \le x_1 \le t+r \\ r & \text{for } x_1 \ge t+r. \end{cases}$$

(4 points)

(f) Recover \bar{u} from \tilde{U} using a certain property of spherical means. (3 points)

Observe that \bar{u} does not depend on x_3 . So by part (b) we have a solution to the 2-dimensional wave equation:

$$u(x_1, x_2, t) = \begin{cases} 0 & \text{for } x_1 < -t \\ 0.25 & \text{for } x_1 = -t \\ 0.5 & \text{for } -t < x_1 < t \\ 0.75 & \text{for } x_1 = t \\ 1 & \text{for } x_1 > t. \end{cases}$$

This solution has jump discontinuities, but this is unsurprising since the initial conditions also had them.

43. Wave energy modes

Let us work with n = 1 for simplicity. If we apply the Fourier transform to the wave equation, we arrive at

$$\frac{\partial^2 \hat{u}}{\partial t^2} + 4\pi^2 k^2 \hat{u} = 0.$$

This leads us to define the spectral energy density

$$2\mathcal{E}(k,t) := \left| \frac{\partial \hat{u}}{\partial t} \right|^2 + 4\pi^2 k^2 |\hat{u}|^2.$$

- (a) Through calculation, show that \mathcal{E} is constant in t. Note that \hat{u} is a complex valued function and it is important to respect complex conjugation: $|g|^2 = g\bar{g}$. (1 point)
- (b) Consider a complex valued function g and its conjugate $h = \bar{g}$. Prove that $\hat{h}(k) = \overline{\hat{g}(-k)}$.

 (1 point)
- (c) Using an equation on page 73 for the script, establish Plancherel's theorem (1 point)

$$\int_{\mathbb{R}} |h|^2 dx = \int_{\mathbb{R}} |\hat{h}|^2 dk.$$

(d) Use this to show that the energy of a wave is constant. (2 points)