

b) $\Delta h = \text{const}$

Generalise

Ansatz $u(x,t) = \sum_{k=0}^{\infty} h_k(x) t^k$

$\partial_t u = \sum_{k=0}^{\infty} k h_k t^{k-1} = \sum_{k=0}^{\infty} (k+1) h_{k+1} t^k$

$\Delta u = \sum_{k=0}^{\infty} (\Delta h_k) t^k$

$0 = \partial_t u - \Delta u = \sum_{k=0}^{\infty} [(k+1)h_{k+1}(x) - \Delta h_k(x)] t^k$

$\Rightarrow h_{k+1}(x) = \frac{1}{k+1} \Delta h_k(x)$

$h_k(x) = \frac{1}{k!} \Delta^k h_0$

$u(x,t) = \sum_{k=0}^{\infty} \frac{1}{k!} (\Delta^k h) t^k$

$u(x,0) = h_0(x) + 0 + 0 + \dots$

If $\Delta h = c \quad \Delta c = 0$

$u = h + \frac{1}{1!} c t + \frac{1}{2!} c^2 t^2 + \dots$

(d) $v(x,t) = u(\lambda x, \lambda^2 t)$ is also a soln

$\lambda=2$
 $v(0,t)$ vs $v(1,t)$ vs $u(0,4t)$ vs $u(2,4t)$



Consider $u(x,t) = e^{-t} \sin x$
 $v(x,t) = e^{-4t} \sin(2x)$

e^{ix}

$e^{2i\lambda k x}$

$e^{-4t^2 k^2 t}$

$\Phi(x, s+t) = \int_{\mathbb{R}} \Phi(y-x, s) \Phi(y, t) dy$

$\Phi(x,t) = \frac{1}{\sqrt{4\pi t}} \exp(-\frac{x^2}{4t})$

Integral formula soln of Cauchy Problem

$u(x,t) = \int_{\mathbb{R}} \Phi(y-x, t) h(y) dy$

$= \Phi(\cdot, t) * h$

$u(x, s+t) = \Phi(\cdot, s+t) * h$

Consider this as the starting temperature of a new Cauchy problem

$u(x, s+t) = \Phi(\cdot, s) * u(\cdot, t)$

$= \Phi(\cdot, s) * [\Phi(\cdot, t) * h]$

$= [\Phi(\cdot, s) * \Phi(\cdot, t)] * h$

This is a dynamical system with evolution operator $\Phi(\cdot, t) *$

$I(0) = I = \int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi}$

$I(k) = \int_{\mathbb{R}+ik} e^{-z^2} dz$ $y = z+ik$ $dy = dz$ $y \in \mathbb{R}+ik \Leftrightarrow z \in \mathbb{R}$

$= \int_{\mathbb{R}} e^{-(z+ik)^2} dz$

$\frac{\partial I}{\partial k} = \int_{\mathbb{R}} \frac{\partial}{\partial k} e^{-(z+ik)^2} dz =$

$= \int_{\mathbb{R}} -2(z+ik) i e^{-(z+ik)^2} dz$

$= i \int_{\mathbb{R}} -2(z+ik) e^{-(z+ik)^2} dz$

$= i \int_{\mathbb{R}} \frac{\partial}{\partial z} [e^{-(z+ik)^2}] dz$

$= i [e^{-(z+ik)^2}]_{z=-\infty}^{z=\infty} = 0$

Lemma

$p(x)$ polynomial

$\lim_{x \rightarrow \pm\infty} p(x) e^{-x^2} = 0$

Proof $\lim_{x \rightarrow \infty} \frac{p(x)}{e^{x^2}} = \lim_{x \rightarrow \infty} \frac{p'(x)}{2x e^{x^2}} = \lim_{x \rightarrow \infty} \frac{p''(x)}{2x^2 e^{x^2}}$

$= \dots = \lim_{x \rightarrow \infty} \frac{C}{e^{x^2}} = 0$

To show that $\exp(-x^2)$ is Schwartz

$\forall \alpha \in \mathbb{N} \sup_{x \in \mathbb{R}} |x^\alpha e^{-x^2}| < \infty$

$f = e^{-x^2}$

$f' = -2x e^{-x^2}$

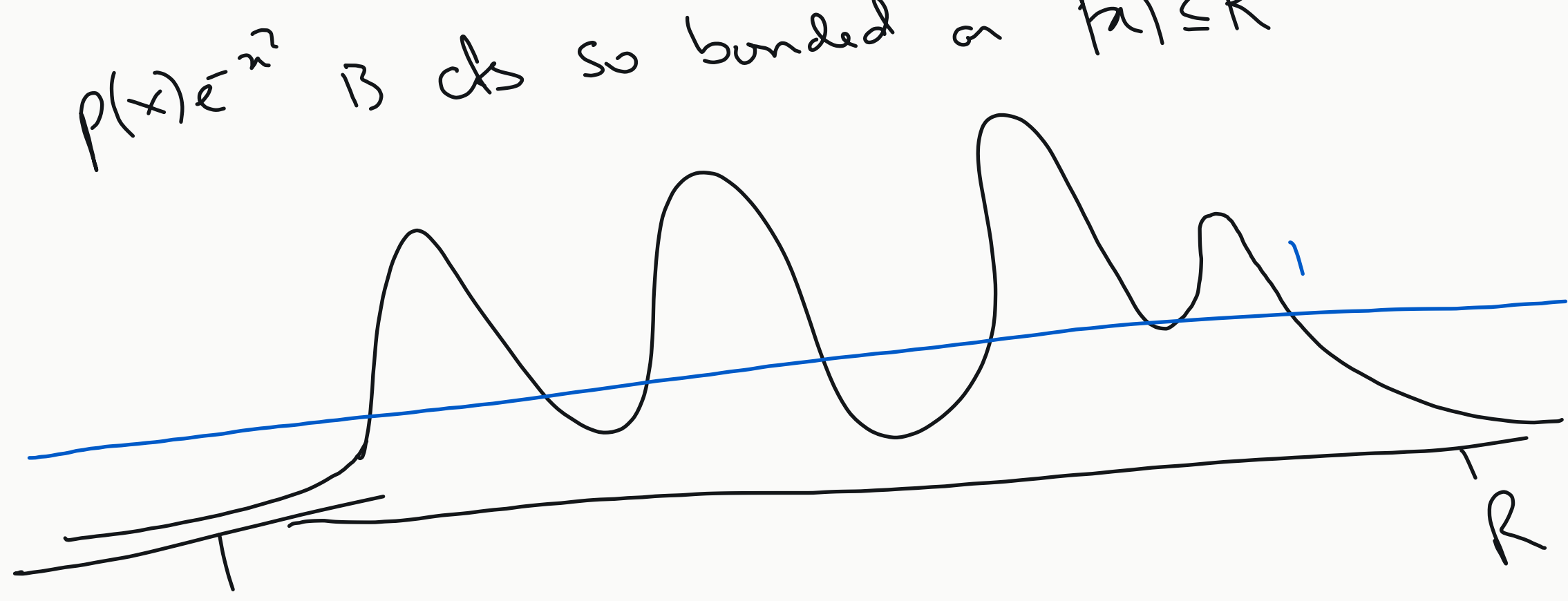
$f'' = -2e^{-x^2} + 4x^2 e^{-x^2} = (2+4x^2) e^{-x^2}$

\vdots
 $p(x) e^{-x^2}$

Need to prove $\sup_{\mathbb{R}} |p(x) e^{-x^2}| < \infty$

$\varepsilon = 1 \exists R \forall |x| > R \quad |p(x) e^{-x^2}| < \varepsilon = 1$

$p(x) e^{-x^2}$ is obs so bounded on $|x| \leq R$



rescaling $\hat{F}[g(ax)] = |a|^{-n} \hat{g}(a^{-1}k)$

$F[u+v] = \hat{u} \hat{v} \quad F[uv] = \hat{u} \hat{a} \hat{v}$

$\hat{F}[g(a^{-1}x)] = |a|^n \hat{g}(ak)$

$\hat{F}(\phi) = F(\hat{\phi})$

$\hat{\delta}(\phi) = \delta(\hat{\phi}) = \hat{\phi}(0)$

$= \int_{\mathbb{R}} \phi(x) e^{-2\pi i \cdot 0 \cdot x} dx$

$= \int_{\mathbb{R}} \phi(x) dx$

$= F_1(\phi)$

$\hat{\delta} = F_1$ Maximum correct

" $\hat{\delta} = 1$ " only one possible meaning but distributions are not functions

$\hat{1}(k) = \int_{\mathbb{R}} e^{-2\pi i k \cdot x} dx$

not integrable

$= \begin{cases} k \neq 0 & \text{"morally zero"} \\ k = 0 & \int_{\mathbb{R}} 1 dx = \infty \end{cases}$

Not defined, but similar to delta dist δ .