

35. Sugar, we're going down swinging.

First let $\Omega' \subset \mathbb{R}^n \times \mathbb{R}$ be an open and connected region. A function $v : \Omega' \rightarrow \mathbb{R}$ is called a *sub-solution* of the heat equation if $\dot{v} - \Delta v \leq 0$.

- (a) *Mean value estimate for sub-solutions* Take any point $(x, t) \in \Omega'$ and a small radius $r > 0$ so that $E(x, t, r) \in \Omega'$ (refer to Definition 4.6). Modify the proof the mean value property of the heat equation to show that

$$v(x, t) \leq \frac{1}{4r^n} \int_{E(x, t, r)} v(y, s) \frac{|x - y|^2}{|t - s|^2} d^n y ds$$

holds for all sub-solutions. (2 Points)

Now let $\Omega \subset \mathbb{R}^n$ be an open, bounded, and path connected region. We denote the parabolic cylinder of Ω by $\Omega_T := \Omega \times (0, T]$ as in Section 4.4. Suppose that $v : \Omega_T \rightarrow \mathbb{R}$ is a sub-solution that extends continuously to $\overline{\Omega_T}$.

- (b) *Maximum principle for sub-solutions* Following on from (a), establish that if v takes the value $\sup_{\Omega_T} v$ on Ω_T , then it is constant. (2 Points)
- (c) *A monotonicity property* For $j \in \{1, 2\}$ let $f_j : \Omega \times (0, T) \rightarrow \mathbb{R}$, $h_j : \Omega \rightarrow \mathbb{R}$, and $g_j : \partial\Omega \times [0, T]$ be smooth functions, and likewise let $u_j : \Omega \times (0, T)$ be smooth functions with continuous extensions to the boundary that satisfy

$$\begin{cases} \dot{u}_j - \Delta u_j = f_j & \text{on } \Omega \times (0, T) \\ u_j(x, 0) = h_j(x) & \text{on } \Omega \\ u_j = g_j & \text{on } \partial\Omega \times [0, T]. \end{cases}$$

Suppose further that $f_1 \leq f_2$, $g_1 \leq g_2$, and $h_1 \leq h_2$. Show in this case that $u_1 \leq u_2$ as well. (2 Points)

36. Heat death of the universe.

First a corollary to Theorem 4.3:

- (a) Suppose that $h \in C_b(\mathbb{R}^n) \cap L^1(\mathbb{R}^n)$ and u is defined as in Theorem 4.3. Show

$$\sup_{x \in \mathbb{R}^n} |u(x, t)| \leq \frac{1}{(4\pi t)^{n/2}} \|h\|_{L^1}.$$

(2 Points)

The above corollary shows how solutions to the heat equation on $\mathbb{R}^n \times \mathbb{R}^+$ with such initial conditions behave: they tend to zero as $t \rightarrow \infty$. Physically this is because if $h \in L^1$ then there is a finite amount of total heat, which over time becomes evenly spread across the plane.

On open and bounded domains $\Omega \subset \mathbb{R}^n$ we can have different behaviour, due to the boundary conditions holding the temperature steady. In this question we determine the long time behaviour of solutions u to the heat equation on open and bounded sets Ω with $u(x, t) = g(x)$ on $\partial\Omega \times \mathbb{R}^+$ and $u(x, 0) = h(x)$. Assume that there is a steady state solution, i.e. a solution to the Dirichlet problem for the Laplace equation $\Delta v = 0$ and $v|_{\partial\Omega} = g$. We claim $u \rightarrow v$ as $t \rightarrow \infty$. Let $w(x, t) = u(x, t) - v(x)$. The claim is equivalent to $w \rightarrow 0$.

- (b) What PDE and boundary conditions does w obey? (2 Points)
- (c) Let l_m be the function from Theorem 4.3 that solves heat equation on \mathbb{R}^n with $l_m(x, 0) = mk(x)$ for m a constant and $k : \mathbb{R}^n \rightarrow [0, 1]$ a smooth function of compact support such that $k|_{\Omega} \equiv 1$. Why must k exist? Why does $l_m \rightarrow 0$ as $t \rightarrow \infty$? What boundary conditions on Ω does it obey? (3 Points)
- (d) Use the monotonicity property to show that w tends to zero. (2 Points)
 Hint. Consider $a = \sup_{x \in \Omega} |w(x, 0)|$.

37. The Fourier transform.

Recall that the Fourier transform of a function $h(x) : \mathbb{R}^n \rightarrow \mathbb{R}$ is defined in Section 4.6 to be a function $\hat{h}(k) : \mathbb{R}^n \rightarrow \mathbb{R}$ given by

$$\hat{h}(k) = \int_{\mathbb{R}^n} e^{-2\pi i k \cdot y} h(y) dy.$$

Lemma 4.20 shows that it is well-defined for Schwartz functions.

- (a) Give the definition of a Schwartz function. (1 Point)
- (b) Argue that $f : \mathbb{R} \rightarrow \mathbb{R}$ given by $f(x) = \exp(-x^2)$ is a Schwartz function. (1 Point)
- (c) Show that the Fourier transform of $\exp(-A^2 x^2)$ for a constant $A > 0$ is $\sqrt{\pi} A^{-1} \exp(-\pi^2 k^2 A^{-2})$. You may use that $\int_{\mathbb{R}} \exp(-x^2) dx = \sqrt{\pi}$. (2 Points)
- (d) Show that $\widehat{\partial_j f}(k) = 2\pi i k_j \hat{f}(k)$ for Schwartz functions $f : \mathbb{R}^n \rightarrow \mathbb{R}$. (2 Points)
- (e) If $u : \mathbb{R} \times \mathbb{R}^+$ is a solution to the heat equation, we can apply a Fourier transform in the space coordinate to get a function $\hat{u}(k, t)$. Show that this function obeys

$$\frac{\partial \hat{u}}{\partial t} + 4\pi^2 k^2 \hat{u} = 0.$$

Solve this ODE in the time variable. (2 Points)

- (f) Suppose that we have the initial condition $u(x, 0) = h(x)$ for $x \in \mathbb{R}$ for a Schwartz function h . Then $\hat{u}(k, 0) = \hat{h}(k)$. Apply the inverse Fourier transformation to rederive the solution given in Theorem 4.3. (2 Points)