

all partial derivative

30a) Schwartz function decays faster than polynomials.

$$f(x) = e^{-x^2} \rightarrow 0 \text{ for all } n, \text{ for } |x| \rightarrow \infty$$

$$f^{(n)}(x) = p_n(x) e^{-x^2} \rightarrow 0 \text{ for } |x| \rightarrow \infty.$$

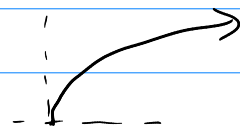
All functions of compact support. Fourier transforms defined  $L^2$

b)  $f(x) = e^{-x^2} (e^{-x^2} + \sin^2 x)$

This is Schwartz basically because  $e^{-x^2}$  term.

$$g = \sqrt{f} = e^{-\frac{1}{2}x^2} \sqrt{e^{-x^2} + \sin^2 x}$$

Behaviour as  $x \rightarrow \infty$



Observe  $\sqrt{x}$  has  $\infty$  derivative at  $x=0$ .

$$h(x) = e^{-\frac{1}{2}x^2} \sqrt{e^{-x^2} + \sin^2 x} \quad \text{varies at } x=k\pi$$

small

$h(k\pi)$  will have large derivative as  $k \rightarrow \infty$ .

c) For fixed  $t$   $\Phi(x,t) = A e^{-Bx^2} = A e^{-B|x|^2}, x \in \mathbb{R}^n$

Martin writes  $x^2 = |x|^2$ ,  $x \cdot t$  for dot product

Basic property of Fourier series  
on  $S' = [0, 1)$

$$e^{ix}, \sin x, \underbrace{e^{2\pi i x}}_{k, l \text{ integers.}}$$

$$\begin{aligned} \widehat{(e^{2\pi i l x})}(k) &= \int_{S'} e^{-2\pi i k y} e^{2\pi i l y} dy \\ &= \int_0^1 e^{2\pi i y(k-l)} dy \\ &= \begin{cases} \text{if } k \neq l & \frac{1}{2\pi i(k-l)} e^{2\pi i y(k-l)} \Big|_0^1 = 0 \\ \text{if } k = l & y \Big|_0^1 = 1 \end{cases} \end{aligned}$$

$$\begin{aligned} f(x) &= \sum_{l \in \mathbb{Z}} a_l e^{2\pi i l x} = \sum \langle f, e^{2\pi i l y} \rangle e^{2\pi i l x} \\ &= \sum \underbrace{\langle f, e_l \rangle}_{\text{Fourier Series}} e_l \end{aligned}$$

$$\int f(y) e^{-2\pi i k y} dy = a_k$$

$$\langle f, e^{2\pi i k y} \rangle \text{ } L^2\text{-norm.}$$

Fourier transform  $\Rightarrow$  same idea on  $\mathbb{R}$ .

$$f = \int \underbrace{f}_{\hat{f}} e^{2\pi i k x} dk$$

a) In 1-D.  $(x, t) \in \mathbb{R} \times \mathbb{R}^+$   $\Phi = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}}$

$$\begin{aligned} \hat{\Phi}(k, t) &= \int_{\mathbb{R}} e^{-2\pi i k y} \frac{1}{\sqrt{4\pi t}} e^{-\frac{y^2}{4t}} dy \\ &= \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(-\underbrace{2\pi i k y - \frac{y^2}{4t}}_{\substack{a^2 y^2 + 2aby + b^2}}\right) dy \end{aligned}$$

$$\frac{y^2}{4t} + 2\pi i k y + \left(\pi i k \sqrt{4t}\right)^2 - \left(\frac{\pi i k \sqrt{4t}}{-}\right)^2$$

$$= \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp - \left[ \left( \frac{y}{\sqrt{4t}} + \pi i k \sqrt{4t} \right)^2 + \pi^2 k^2 4t \right] dy$$

$$= \frac{\exp(-4\pi^2 k^2 t)}{\sqrt{4t}} \int_{\mathbb{R}} \exp(-z^2) \sqrt{4t} dz$$

$$z = \frac{y}{\sqrt{4t}} + \pi i k \sqrt{4t}$$

$$= \frac{\exp(-4\pi^2 k^2 t)}{\sqrt{4t}} \sqrt{4t}$$

$$e) \sum_{n \in \mathbb{Z}} |f(x+n)| = |f(x)| + |f(x-1)| + \sum_{n \neq 0, -1} |f(x+n)|$$

$x$  is fixed

We know  $f$  decays faster than any polynomial, e.g.  $(x+n)^2$  in  $n$ .

$$\forall \varepsilon > 0 \quad |(x+n)^2 f(x+n)| < \varepsilon \quad \text{as } n \rightarrow \pm \infty$$

$$|f(x+n)| < \frac{\varepsilon}{(x+n)^2}$$

$\exists N$  so that  $\forall |n| \geq N$

choose  $N$  so that  $|x| < N$  for all  $x \in \Omega$

$$\leq \underbrace{\sum_{|n| < N} |f(x+n)|}_{\text{finite}} + \underbrace{\sum_{|n| \geq N} \frac{\varepsilon}{(x+n)^2}}_{\text{convergent}}$$

$$\frac{1}{n^2} \quad x \in \Omega$$

f) If the boundary is periodic show the solution is periodic.

$$(\Delta - \partial_t) u = 0$$

$$u|_{t=0} = h$$

$$\partial_x(u(x+1, t))$$

$$= (\partial_x u)(x+1, t) \times 1$$

$v(x, t) = u(x+1, t) - u(x, t)$  is a sol<sup>n</sup> to heat equation.

$$v(x, 0) = h(x+1) - h(x) = 0 \Rightarrow v(x, t) = 0 \text{ by uniqueness}$$

g)  $S' = [0, 1]$  with  $f(x) = f(x+1)$

Suppose we need a sol<sup>n</sup> to  $(\Delta - \partial_t) u = 0$   
 $u|_{t=0} = h$

is a sol<sup>n</sup> on all of  $\mathbb{R}^n$

Then by representation  $u(x, t) = \int_{\mathbb{R}} h(z) \Phi(x-z, t) dz$

$$= \sum_{n \in \mathbb{Z}} \int_n^{n+1} h(z) \Phi(x-z, t) dz$$

$$z = y+n \quad dz = dy \quad \begin{matrix} z=n & y=0 \\ z=n+1 & y=1 \end{matrix}$$

$$= \sum_{n \in \mathbb{Z}} \int_0^1 h(y) \Phi(x-(y+n), t) dy$$

$$= \int_0^1 h(y) \left[ \sum_{n \in \mathbb{Z}} \Phi(x-(y+n), t) \right] dy$$

heat kernel for  $S'$

32 If you have a function  $f$  on  $[0, 1]$  with  $f(0) = 0 = f(1)$   
 reflect by oddness  $f(-x) = -f(x)$  on  $[-1, 1]$   
 periodic period 2 on  $\mathbb{R}$ .  
 no connection between  $f(0)$  and  $f(1)$

vs.  $S'$   $f$  on  $[0, 1]$   
 periodic with period 1  
 $f(0) = f(1)$   
 $f^{(k)}(0) = f^{(k)}(1)$

$$34/ \quad g(x) = \int_0^\infty \Phi_H(x,t) dt = a \Phi_L(x) + b$$

$$= \left( \int_0^1 + \int_1^\infty \right) \frac{1}{(4\pi t)^{n/2}} e^{-\frac{|x|^2}{4t}} dt.$$

$$\text{as } t \rightarrow 0 \quad -\frac{x^2}{4t} \rightarrow -\infty \quad \Phi_H(x,t) \rightarrow 0 \quad x \neq 0$$

so  $\int_0^1 \Phi_H dt$  is finite

$$\int_1^\infty \Phi_H dt \leq \int_1^\infty \frac{1}{(4\pi t)^{n/2}} = \text{finite } n \geq 3 \quad (\Delta - \partial_t) \Phi_H = 0$$

$$\Delta g = \int_0^\infty \Delta \Phi_H dt = \int_0^\infty \partial_t \Phi_H dt$$

$$= \Phi_H(x, \infty) - \Phi_H(x, 0)$$

$$= 0 - 0$$

$$= 0$$

In general  $\Delta u = \partial_t u$  and apply Laplace transform in  $t$

$$v(x, \lambda) := \int_0^\infty u(x, t) e^{-\lambda t} dt$$

$$\Delta v = \lambda v + u(x, 0)$$

$$\Delta \tilde{\Phi}_H = \lambda \tilde{\Phi}_H + 0$$

For 1D  $(x, t) \in \mathbb{R} \times \mathbb{R}^+$

$$v_{xx} = \lambda v + h(x) \quad \text{an ODE with parameter.}$$